

Terrestrial low-frequency bursts: Escape paths of radio waves through the bow shock

A fascinating investigation

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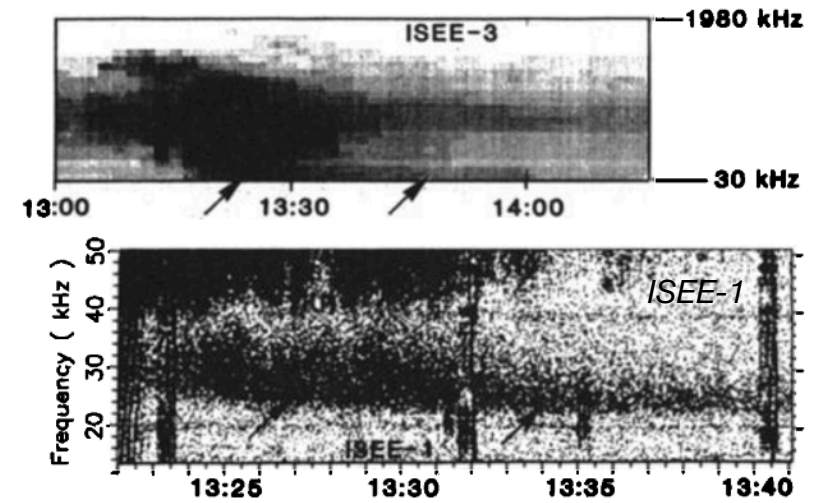
Planetary and Space Science, 52, 643–660, 2004

**and JLS's last paper*

Context

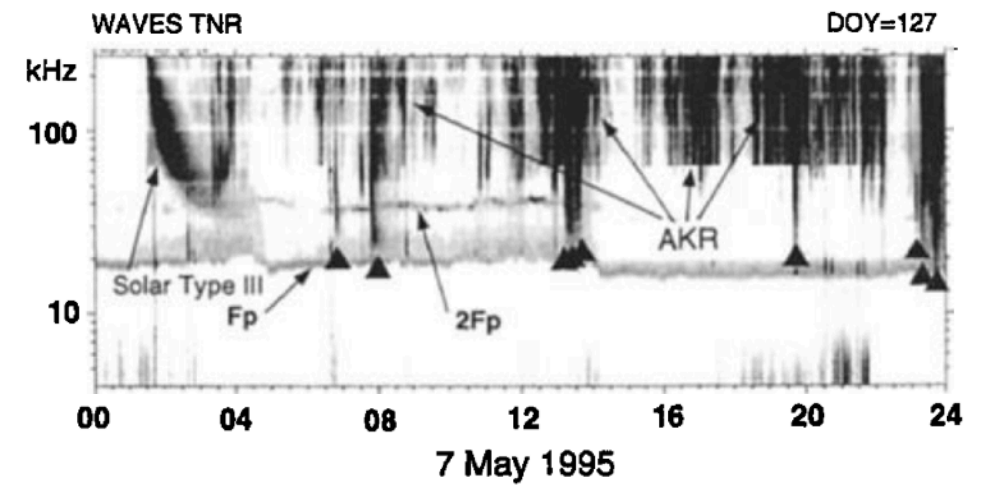
- ITKR : the isotropic terrestrial kilometric radiation

[Steinberg, Lacombe, Hoang, GRL 1988]



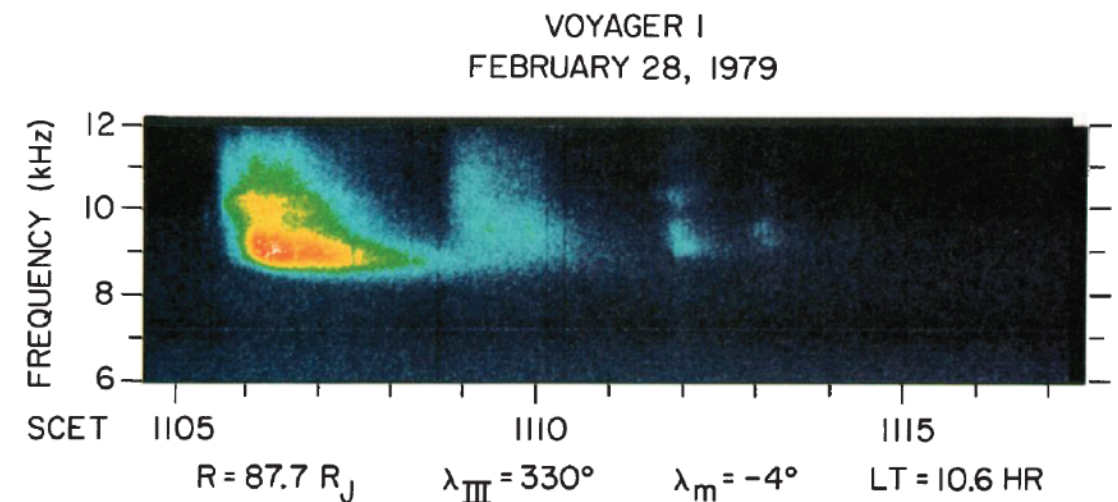
- LF bursts

[Kaiser et al., GRL 1996; Desch et al., GRL 1996]



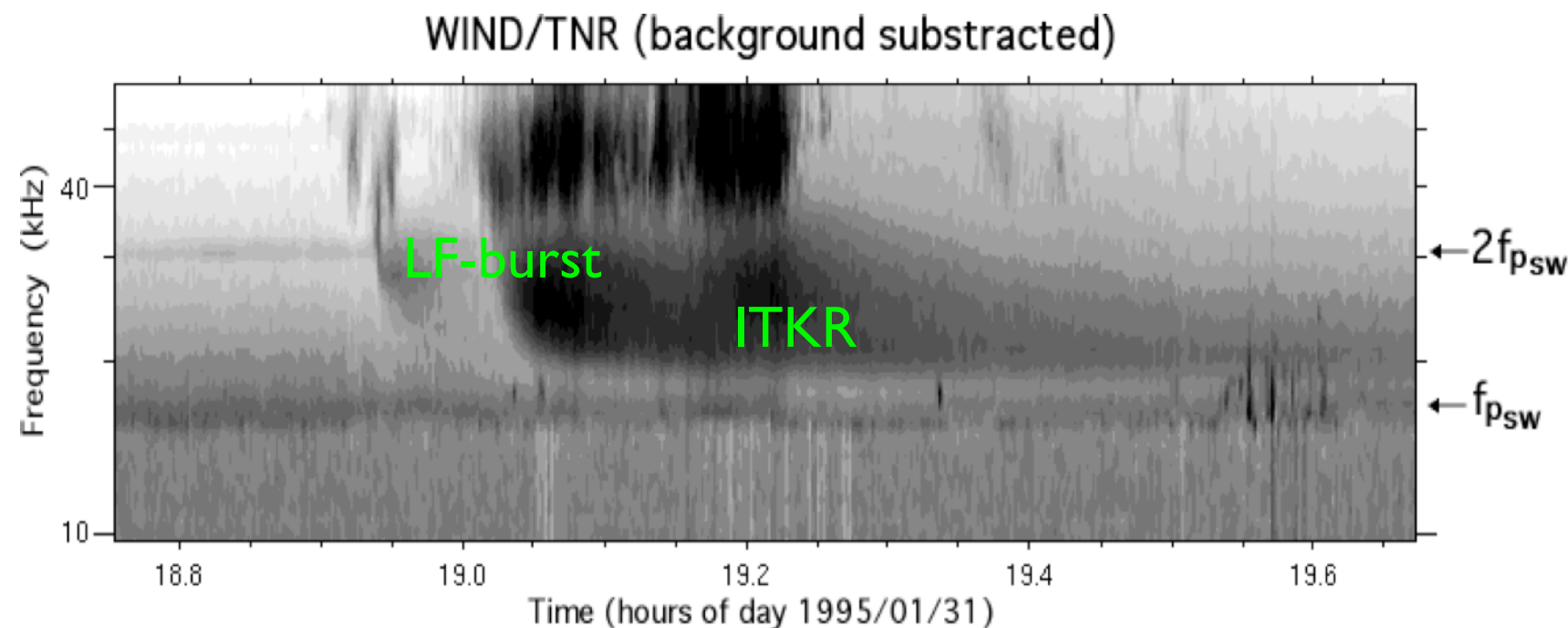
- Jovian type III bursts / QP bursts

[Kurth et al., JGR 1989; MacDowall et al., PSS 1993]



Context

- Correlated with AKR (substorms, SW speed & IMF, auroral source near Earth)
- AKR LF cutoff at $2 f_{psw} = \text{BS nose} \leftrightarrow \text{LF burst} / \text{ITKR} < 2 f_{psw}$
- $P_{\text{LF burst}} \sim 500 \text{ kW} = \text{a few \% of AKR power}$
- LF burst origin : Sporadic electron acceleration/injections ? (reconnection ? tail ?...)
- LF burst: small source (a few $^{\circ}$) [Kaiser et al., GRL 1996]
+ long diffuse tail \sim isotropic (ITKR)



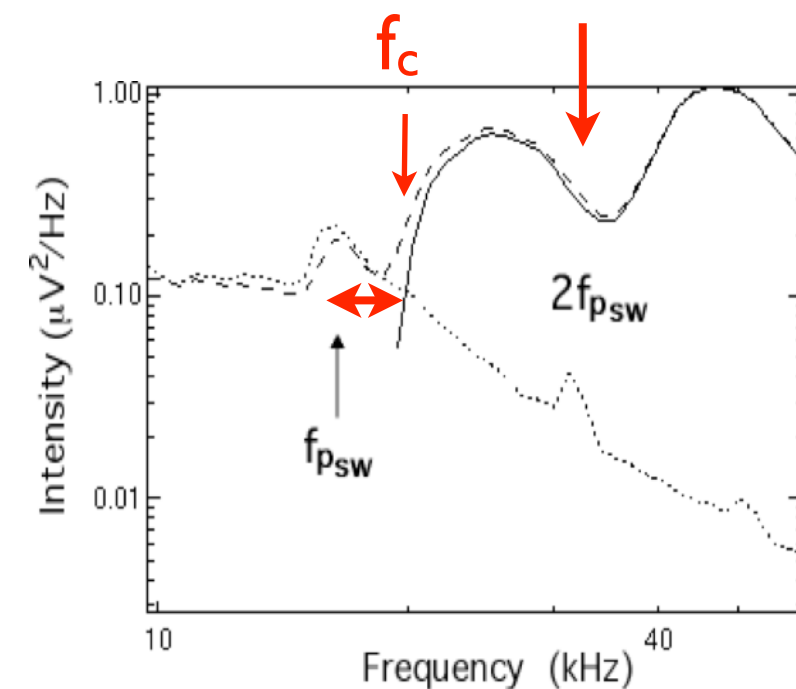
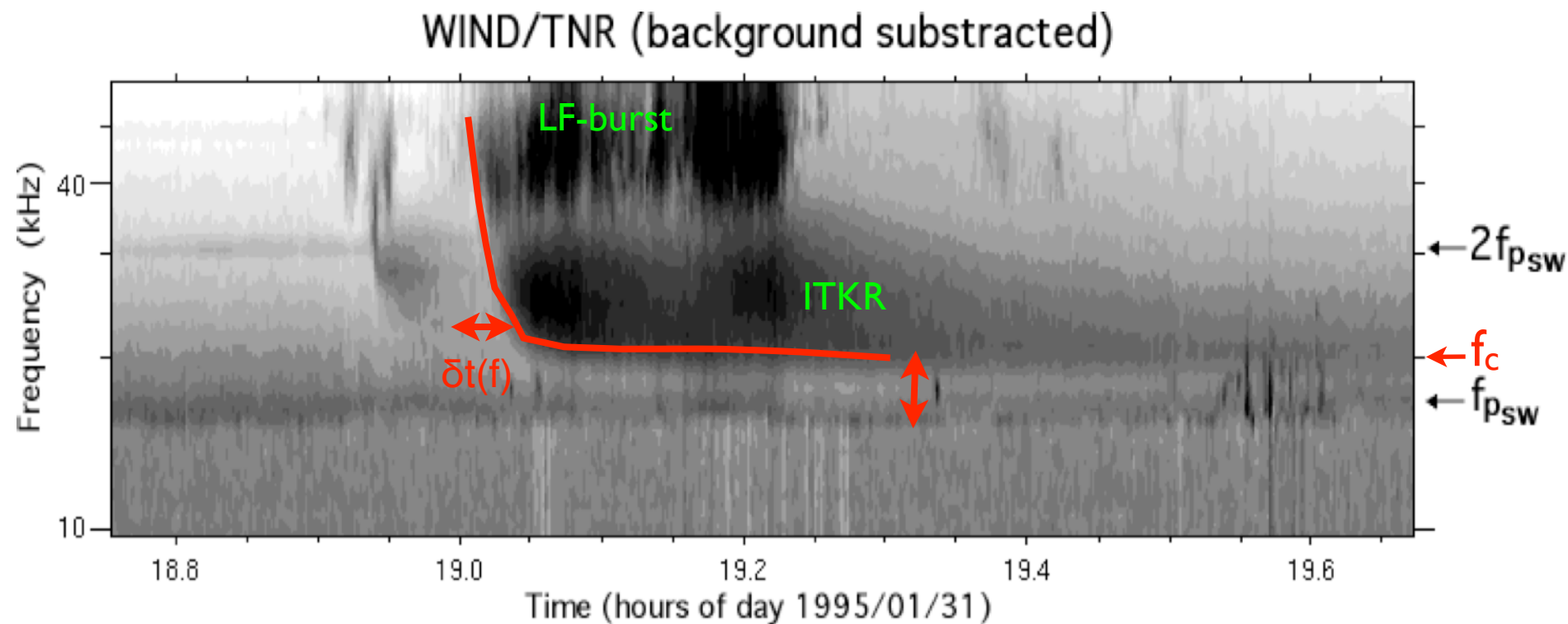
\Rightarrow tailward propagation to/from/through the Mtail ?

in Msheath [Desch & Farrell, 2000] ? but requires $V < V_g$

MS escape far downstream (100's R_E) ?

The Study

- 119 bursts analyzed in Wind/Waves observations near L1 (60-200 R_E upstream of Earth) 12/1994 - 5/1995
- Spectrum from $2 f_{psw}$ down to $f_c > f_{psw}$ (f_c strongly correlated to $f_{psw} \Rightarrow$ propagation-related)

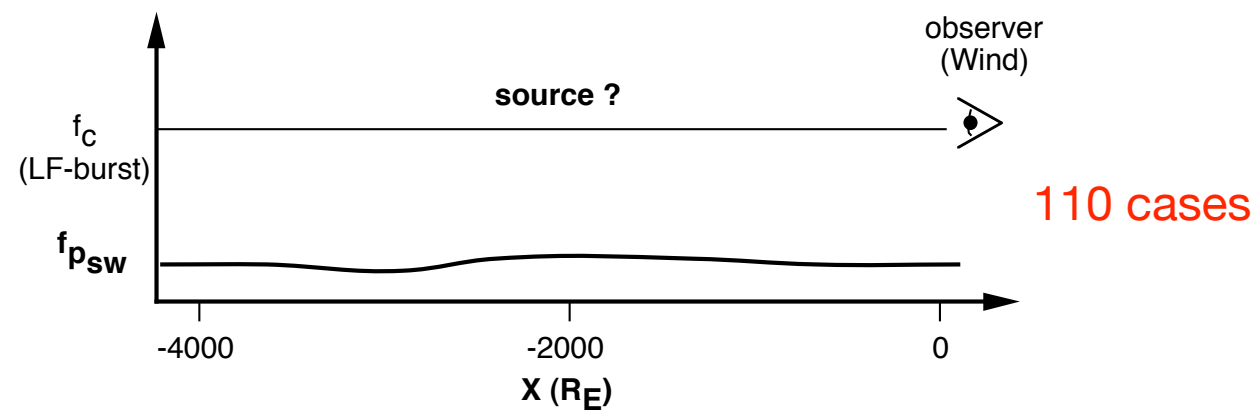
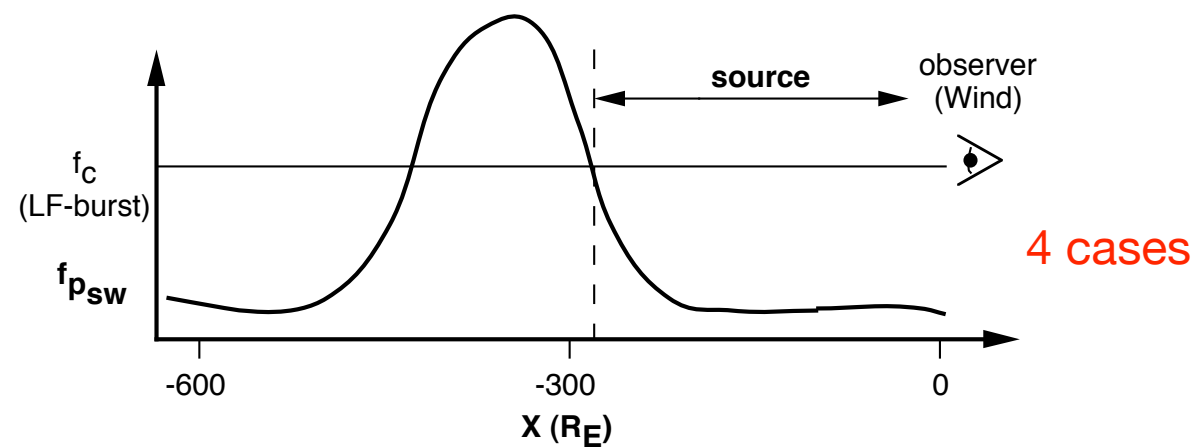
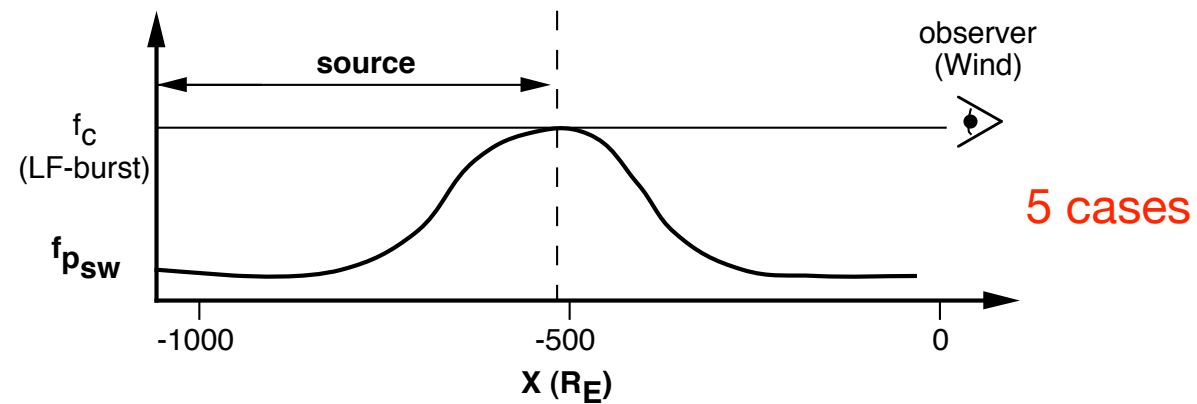


Questions

- Why $f_c > f_{psw}$? (average $\sim 1.28 f_{psw}$)
- $\delta t(f)$? (minutes)
- Spectral dip at $\sim 2 f_{psw}$?

Results (1)

- SW parameters (t) measured every 6 minutes \Rightarrow projected at Earth
 \Rightarrow generally no SW overdensity $\geq f_c$



Model (1)

- **Geometrical model for the Earth's bow shock:**

hyperboloid of revolution

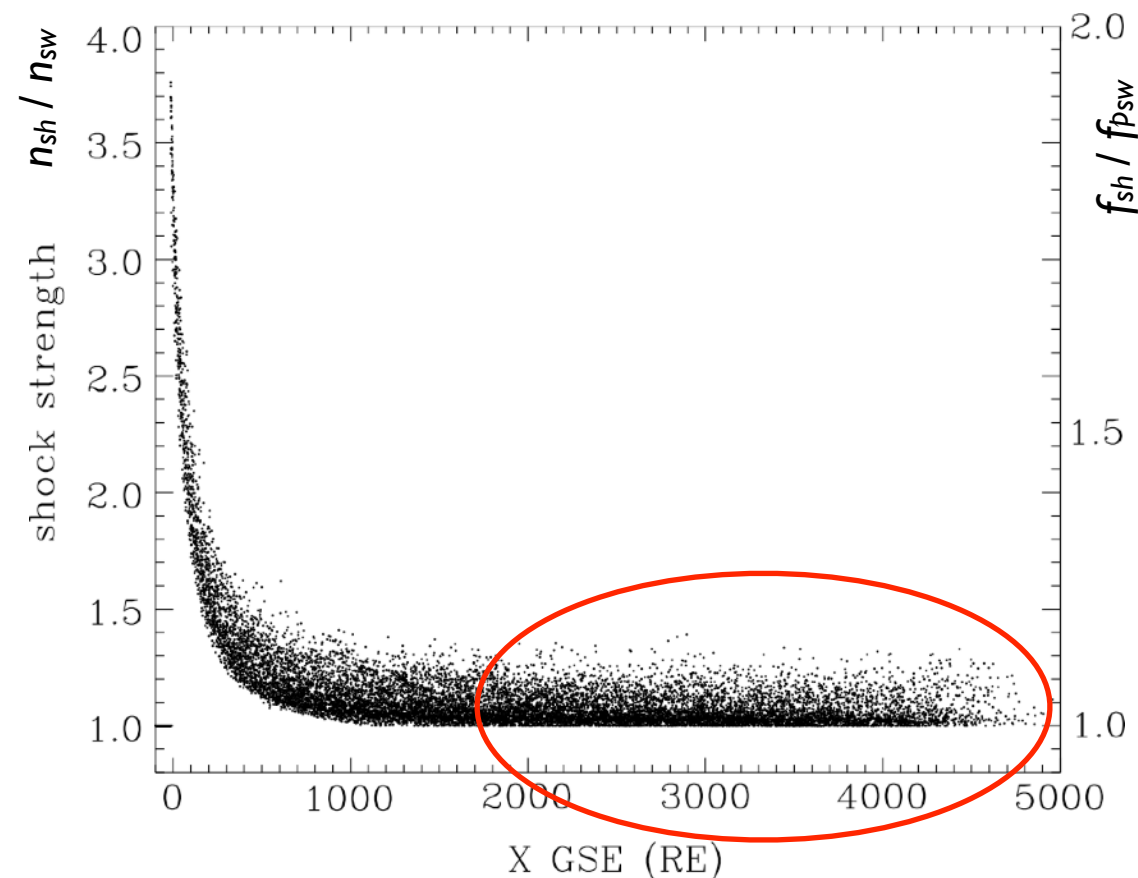
asymptote = fast mode Mach cone of aperture α ($\sin \alpha = 1/M_f$)

+ orbital aberration

⇒ α and orientation depend on SW parameters (N_e , β)

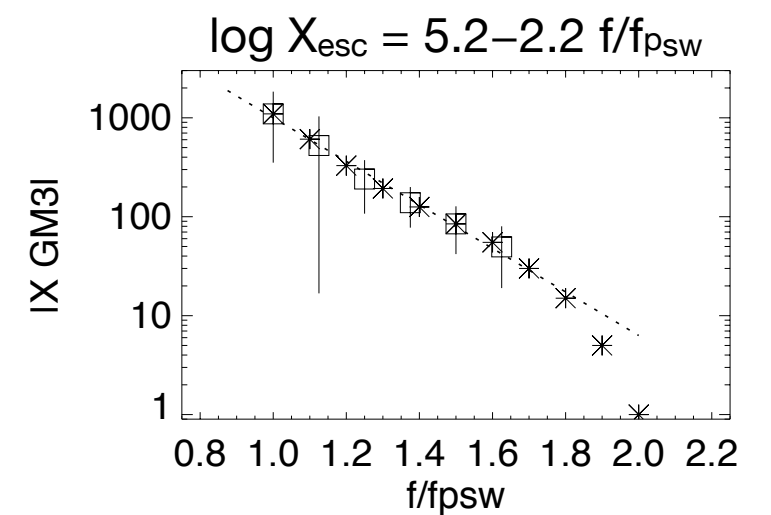
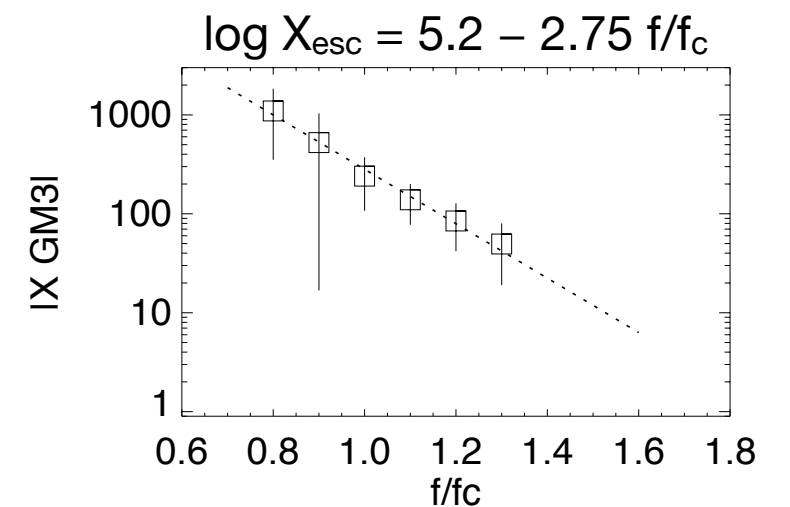
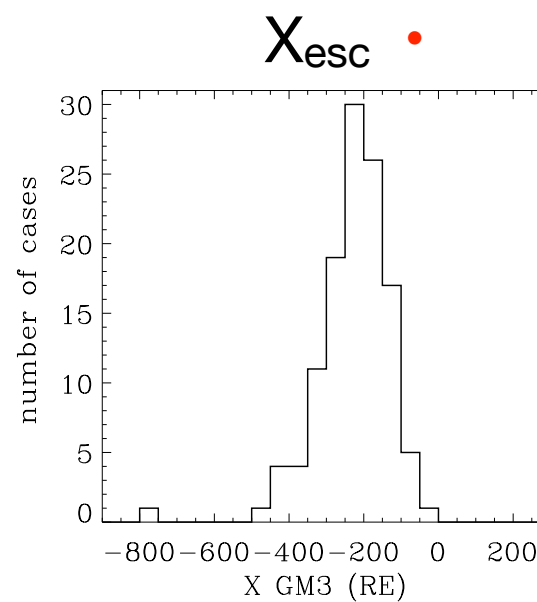
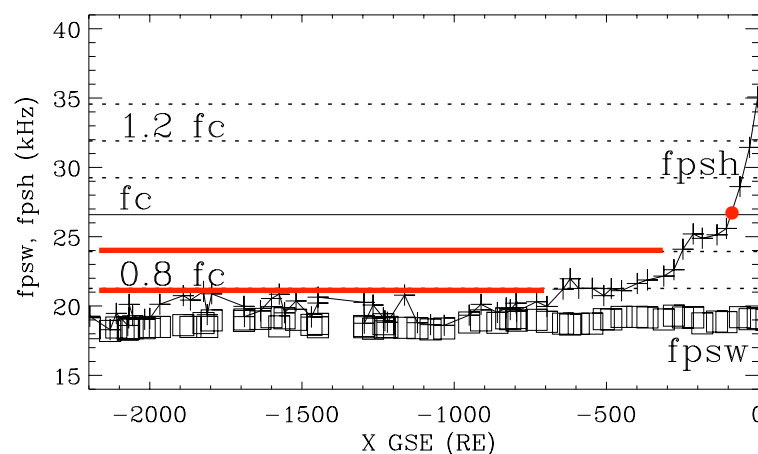
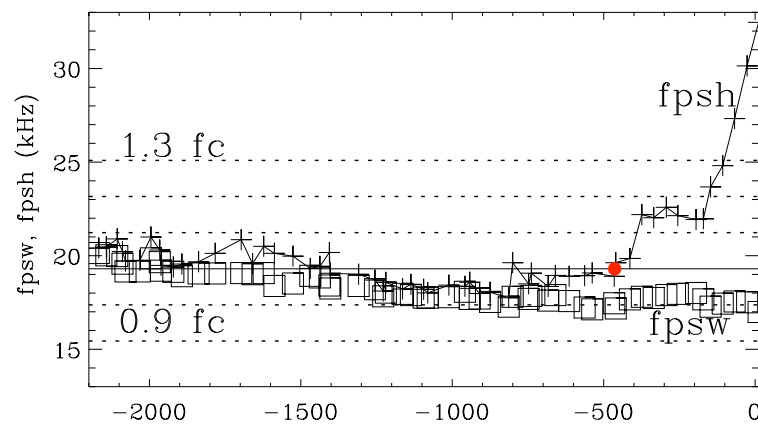
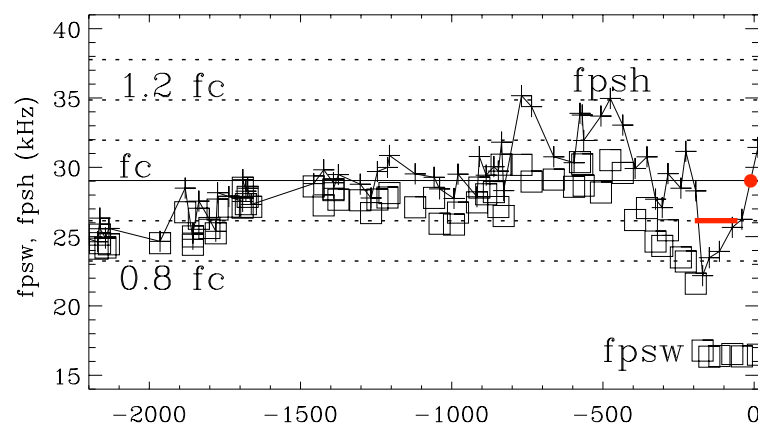
⇒ θ_{BN} vary around & along shock, N_e jump from Rankine-Hugoniot relations

= **GM3** (Greenstadt, Mach, 3D) model



Results (2)

- GM3 BS : shape, orientation, N_e jump for each observed LF burst
 - ⇒ f_{psh} down to 2000-5000 R_E downstream, for each event
 - ⇒ X_{esc} where $f_c \geq f_{\text{psh}}$, ~ 100 -500 R_E , increases exponentially vs. f/f_{psw} and f/f_c
 - ⇒ but this cannot explain the cutoff at f_c rather than e.g. $0.9 f_c$ or $0.8 f_c$



TAIL

WIND

Model (2)

- **Ray-tracing in a scattering homogeneous plasma**

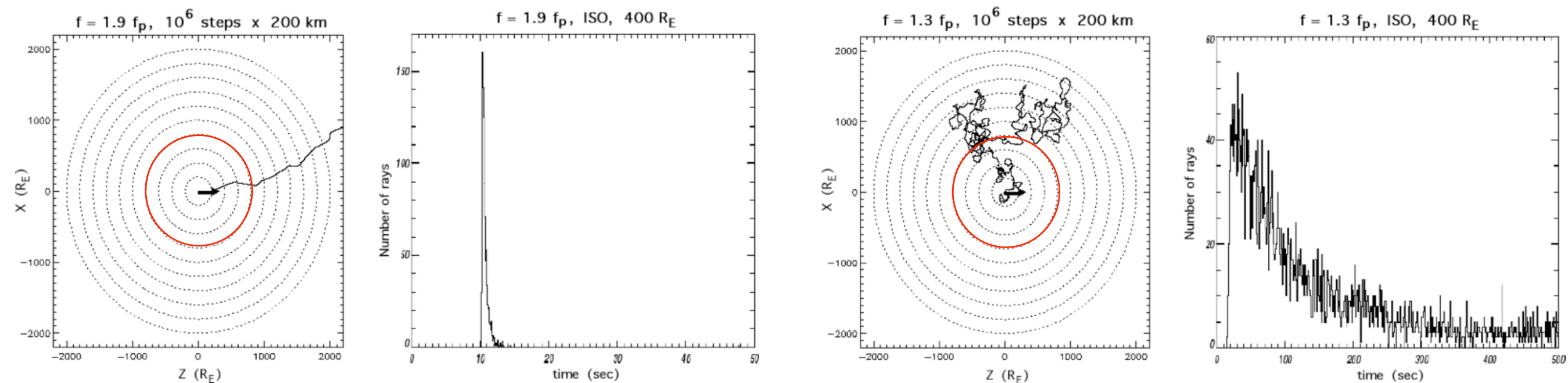
Weak SW N_e gradient \Rightarrow locally homogeneous ($N_e, \delta N_e$)

Scattering by $\Delta\theta^2 = b(f) \Delta S$ (scattering coeff. $b(f)$ [Lacombe et al., AnG 1997], $\Delta S=200$ km

1000 rays launched at $1.1-2 \times f_{p\text{sw}}$, $1-2.5 \cdot 10^6$ steps each = $3-8 \cdot 10^4 R_E$
($\Delta\theta, \phi$) at each step \Rightarrow frame translated + rotated ($\Delta S, \Delta\theta, \phi$)

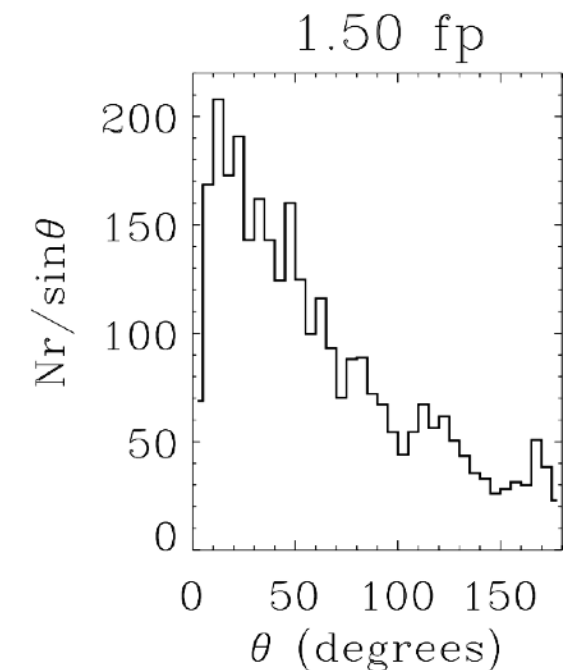
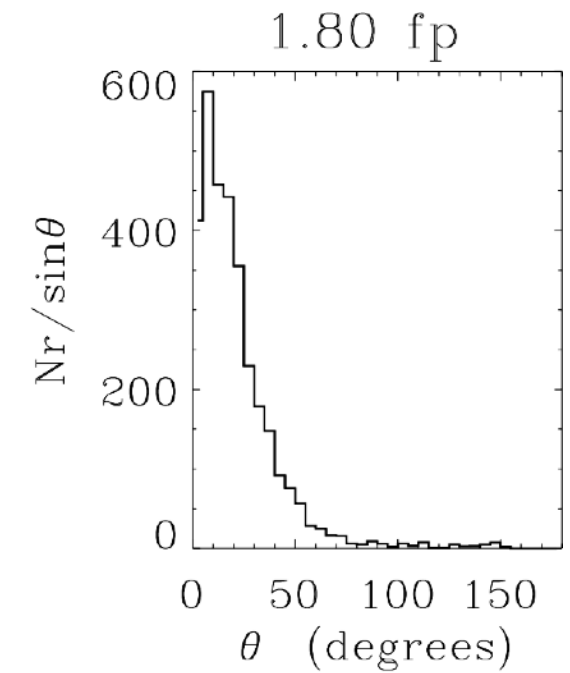
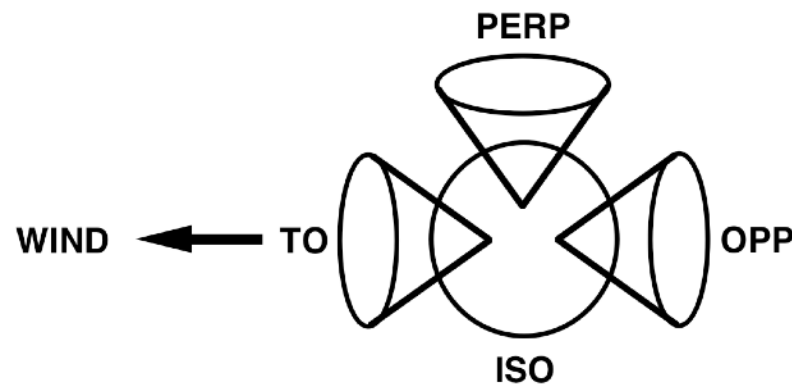
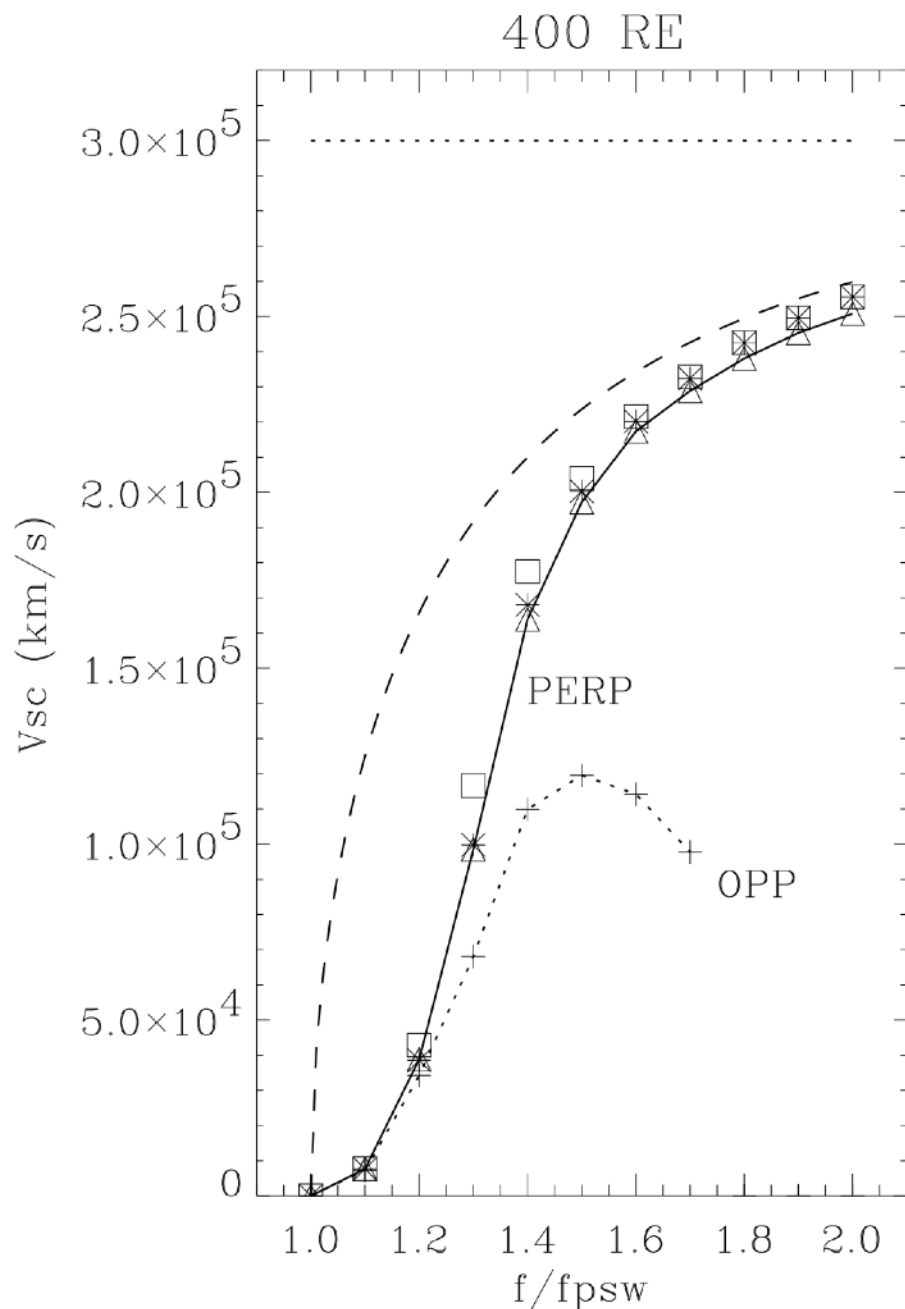
At each step: R , direction(θ, ϕ), time $t=s/V_g$ ($s=k \cdot \Delta S$, $V_g=c\mu$)

N_{rays} counted at 200-2000 R_E vs. frequency (lost at $R \geq 3000 R_E$)



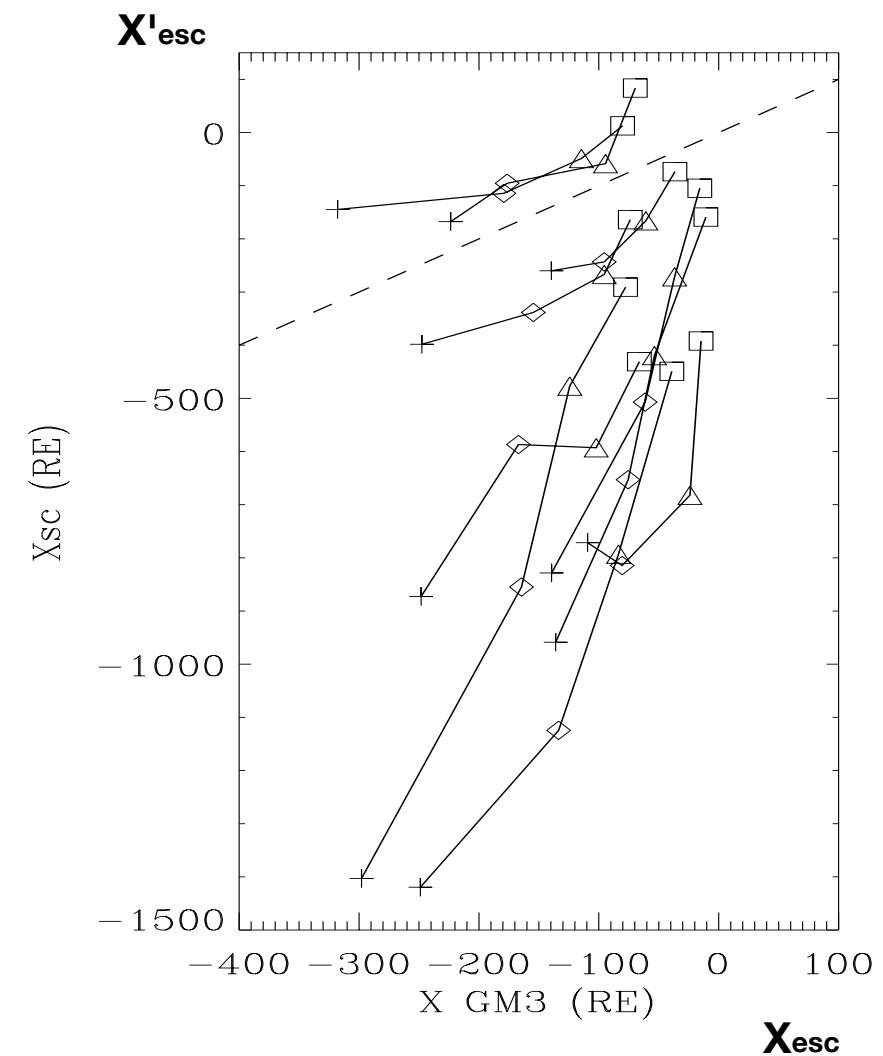
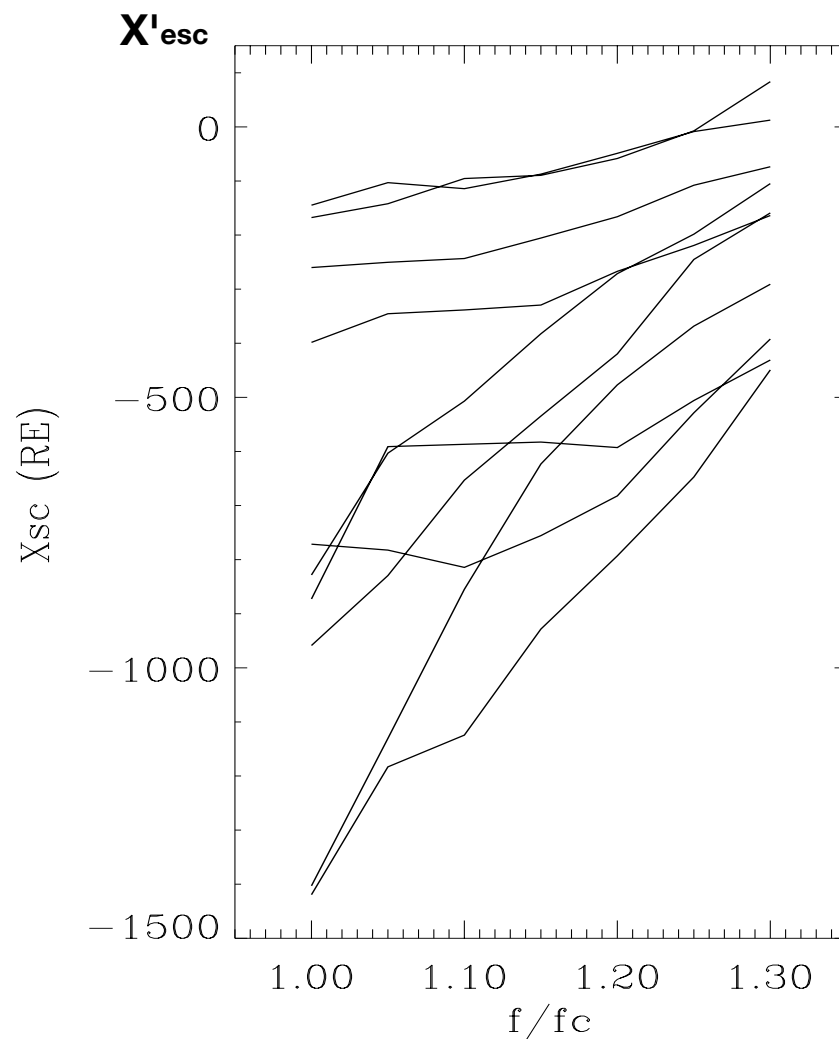
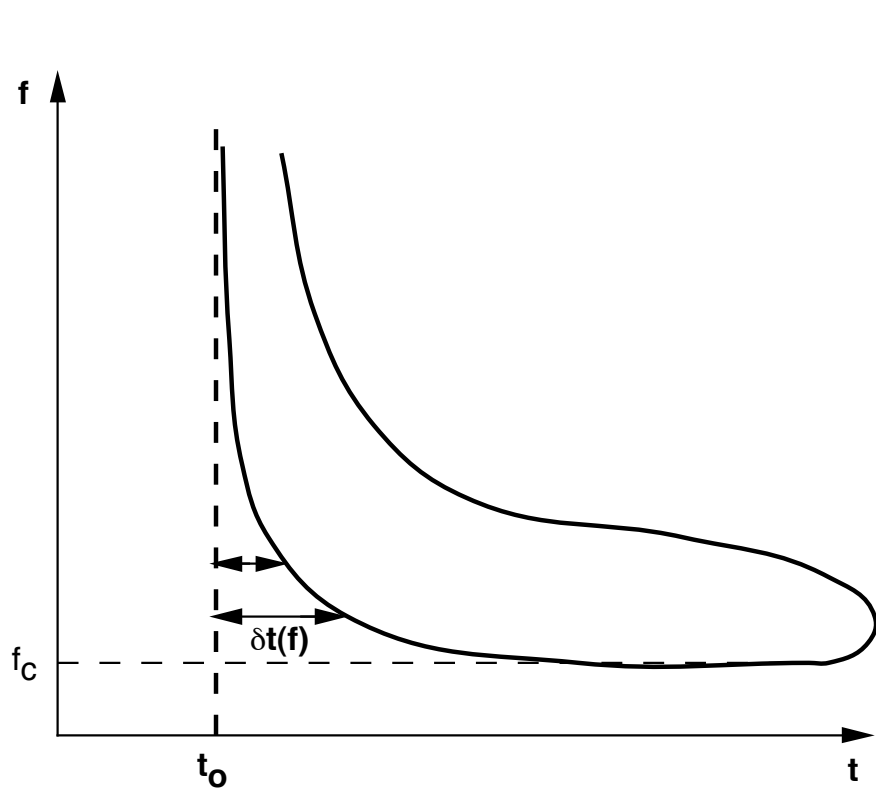
Model (2)

- $V_{sc}(\text{scattered rays}) = t(5\% \text{ rays reach } R)$
⇒ decreases with f/f_{psw} & when R, t increases
- Angular distribution ⇒ isotropisation with decreasing f/f_{psw} & increasing R, t



Results (3)

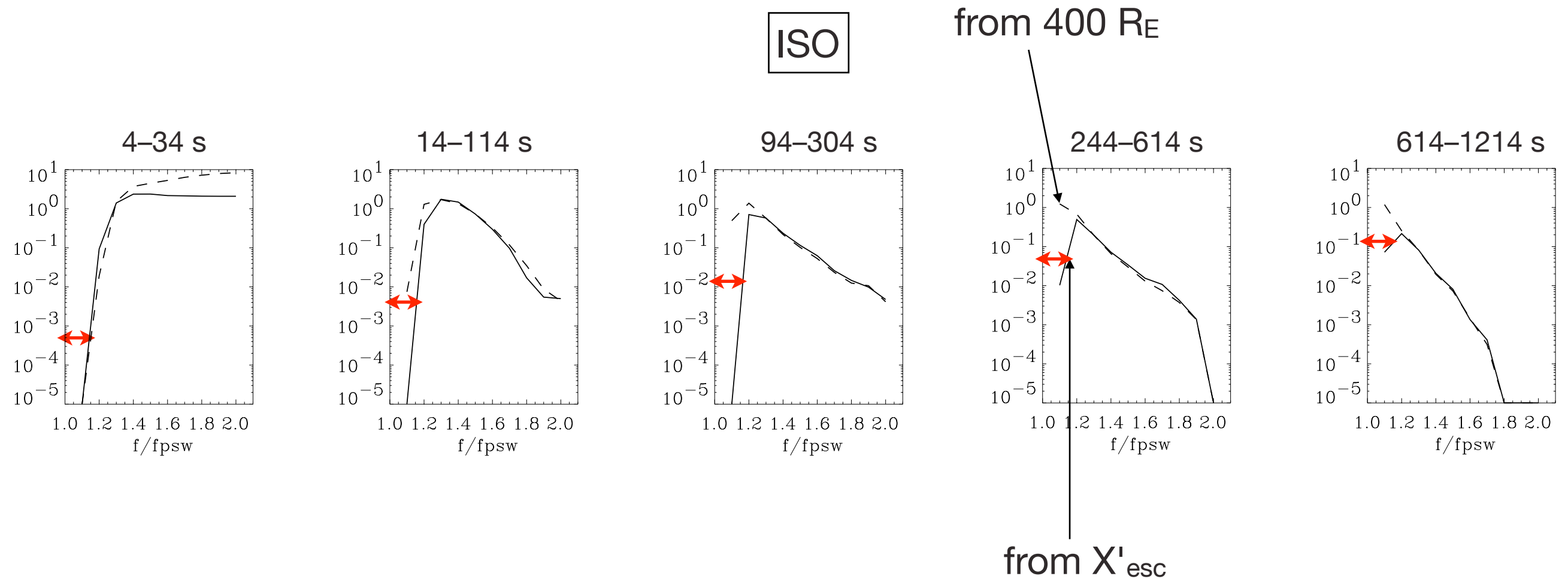
- Scattering in SW or Msheath (locally homogeneous) $\Rightarrow V_{sc}$ decreases with f/f_{psw}
- Propagation through MS(wave guide) + MSheath & SW (scattering, \sim)
 \Rightarrow observed $\delta t(f/f_{psw}) \rightarrow X'_{esc}$ (independent of shock model) $> X_{esc}$
 (same trend \Rightarrow possible to reconcile them via $b(f)\uparrow$, $V_{sc}\downarrow$, GM3 profile ($X_{esc}\uparrow$, N_e jump \uparrow), Msheath guides waves downtail)



Results (3)

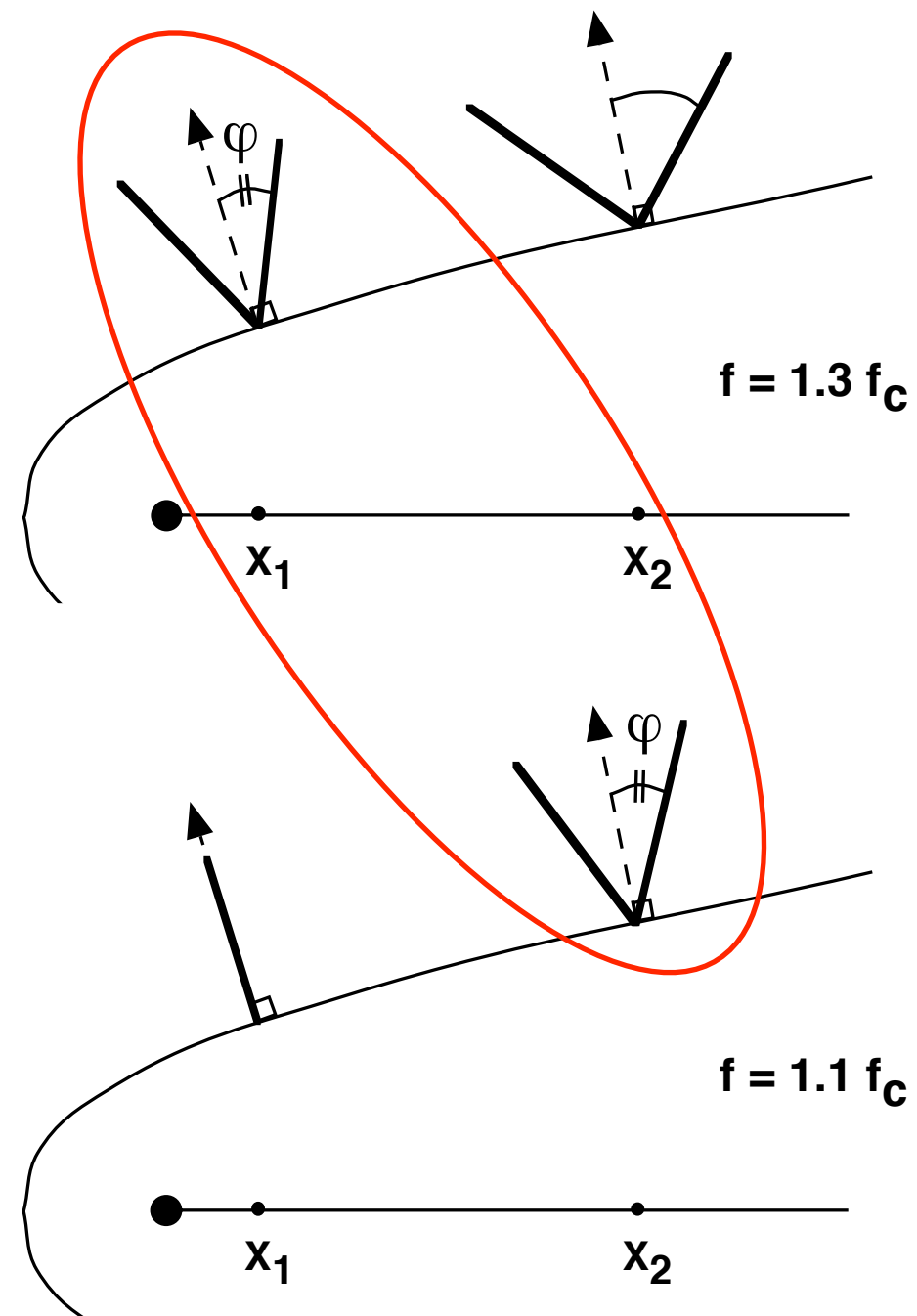
- Exponential increase of X'_{esc} + scattering \Rightarrow explains cutoff at $f_c > f_{psw}$ (rather than $< f_c$)

at e.g. $0.9 f_c \Rightarrow X'_{esc} \uparrow, V_{sc} \downarrow, \delta t \uparrow \Rightarrow$ time dilution + $1/R^2$ intensity decrease \Rightarrow signal undetectable



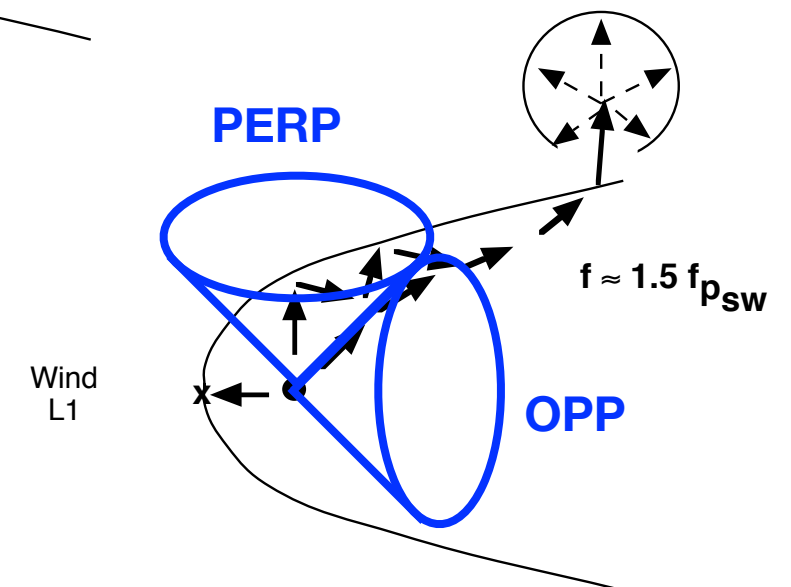
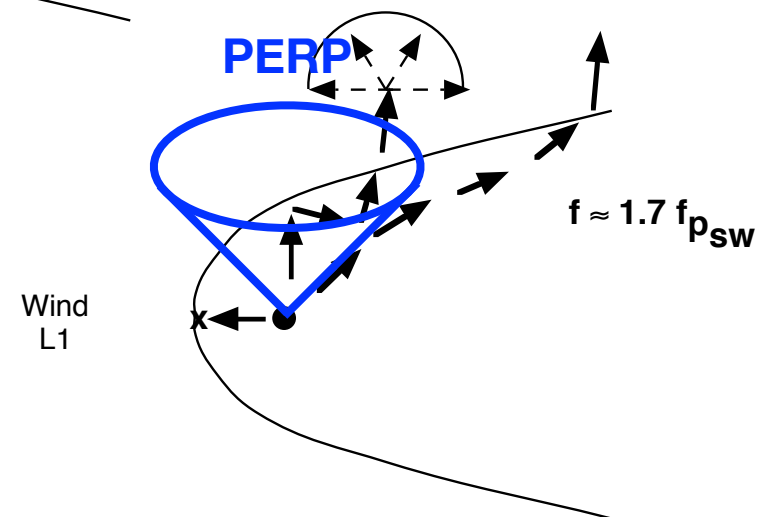
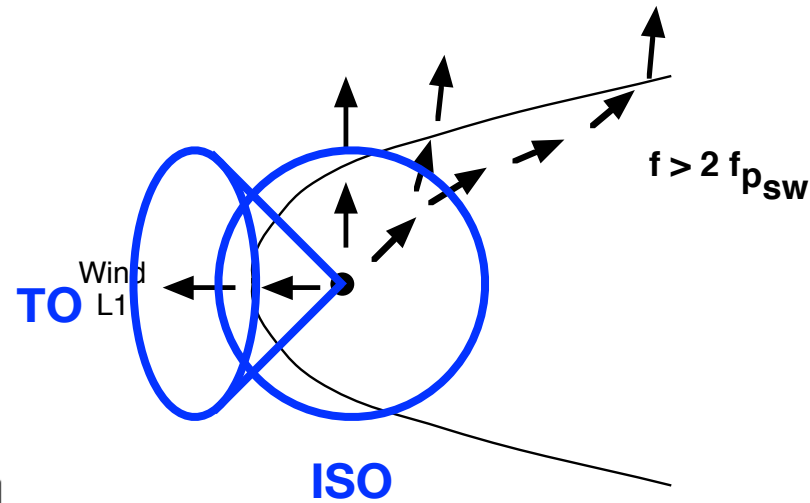
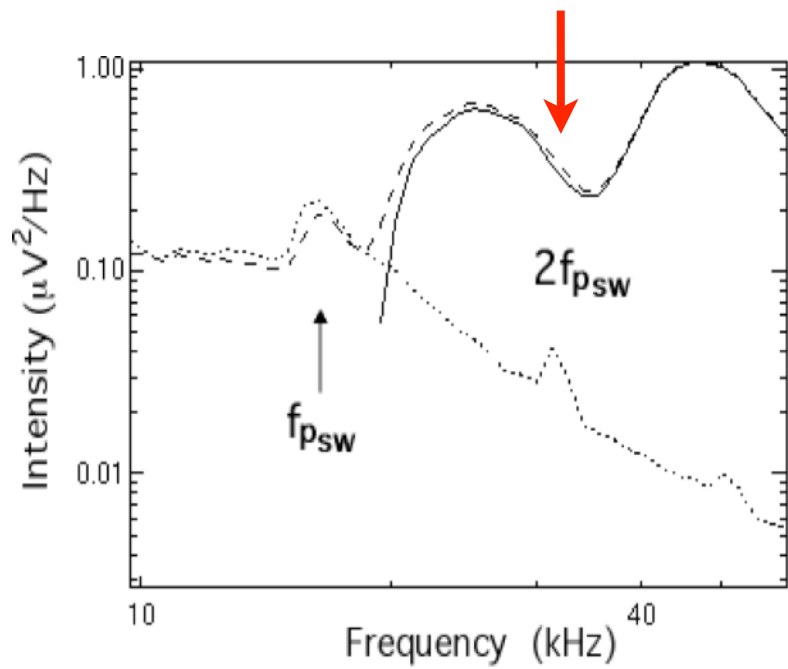
Results (4)

- Density jump at BS : Snell-Descartes law
⇒ exit cone $\sim 20^\circ$ at all f/f_c values because X_{esc} & jump variations compensate



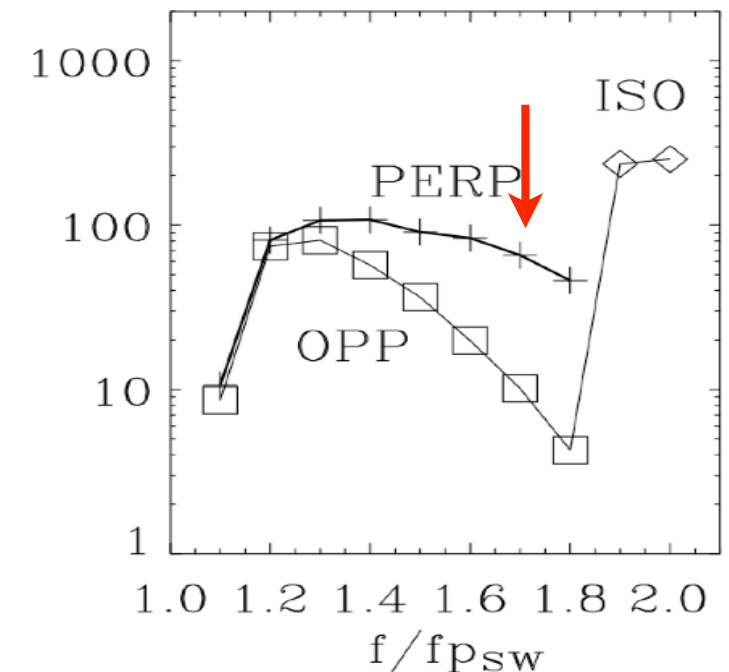
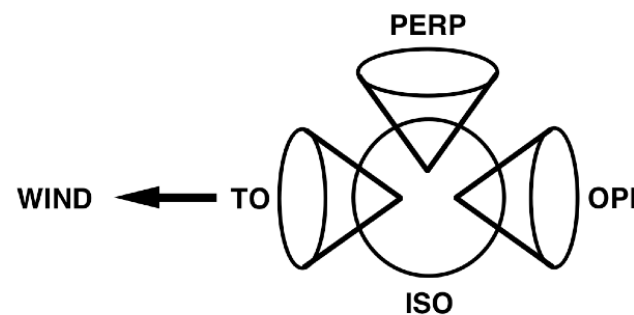
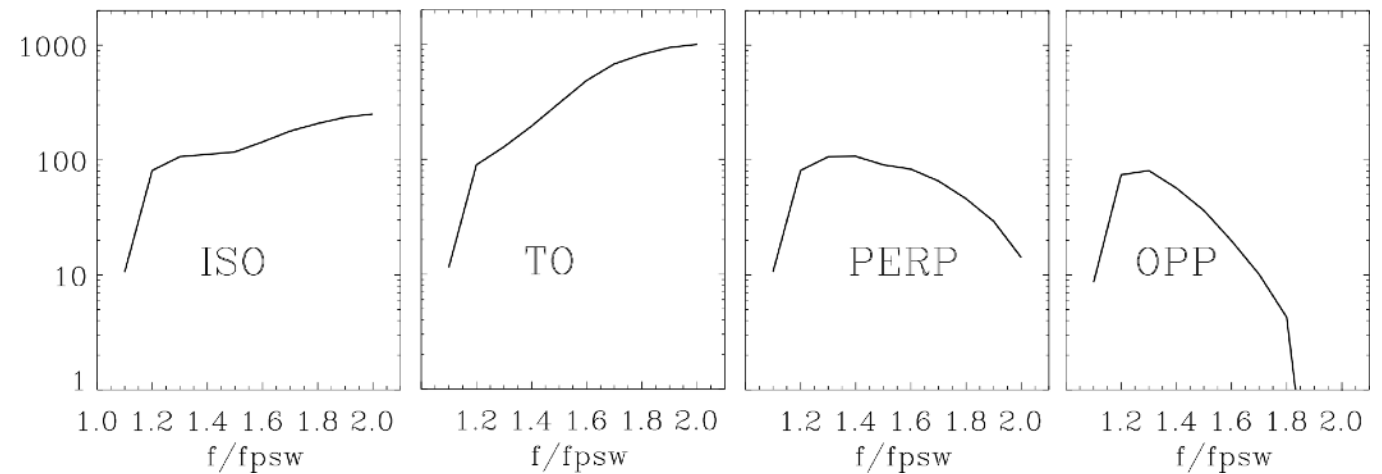
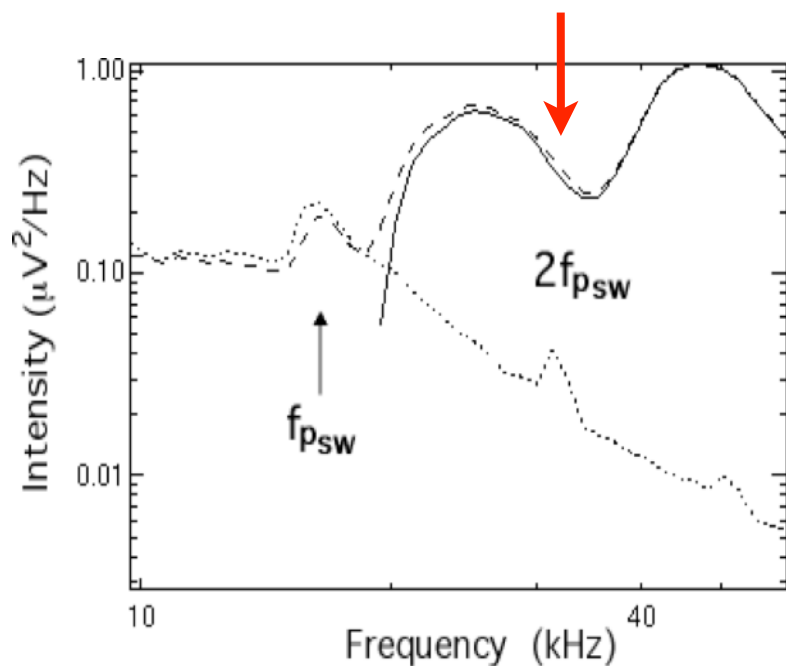
Results (5)

- Two-bump spectrum explained by angular distribution of rays exiting the BS
⇒ Angular distribution broadens with decreasing f and with time



Results (5)

- Two-bump spectrum explained by angular distribution of rays exiting the BS
 - ⇒ Observed spectrum depends on angular distribution at shock exit
 - ⇒ Composite directivity of LF-burst + ITKR explains observed spectrum
 - ⇒ PERP or OPP <math> < 2 f_{psw}</math>, ISO <math> < 2 f_{psw}</math> (no spectral continuity, ≠ scattering)



- Spike seen directive at lower f from BS flanks / dayside ⇒ Ok with model
- At the same frequency, emission can be spiky then diffuse, depending on X'_{esc}

Conclusions

- f_c & $\delta t(f)$ explained by escape distance $X'_{esc} \propto X_{esc} = 10^{5.2-2.2 f/f_{psw}}$ + scattering in SW
- Spectrum explained by angular distribution of rays leaving the shock + scattering

Consequences

- GM3 shock model Ok, locally isotropic scattering Ok
- shock exists $> 1000 R_E$

Open Questions

- Source of LF burst / ITKR ?
- Relation to AKR, to substorms ?
- Why no AKR diffuse tail ?
 - ⇒ HF ⇒ LF burst = LF tail of AKR
 - ⇒ ~? LFE events at Saturn

[Reed et al. JGR 2018; Pazamickas et al. AGU 2004]

Further studies

- LF burst tail $\sim f_{p_{sw}}$ when observed in the MS ⇒ ?

[Anderson et al., AGU 2002, COSPAR 2005]

